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#### A comparative study of manhole hydraulics using 4 stereoscopic PIV and different RANS models 5

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#### 13 Abstract

14 Flows in manholes are complex and may include swirling and recirculation flow with significant 15 turbulence and vorticity. However, how these complex 3D flow patterns could generate different 16 energy losses and so affect flow quantity in the wider sewer network is unknown. In this work, 17 2D3C stereo Particle Image Velocimetry (PIV) measurements are made in a surcharged scaled 18 circular manhole. A CFD model in OpenFOAM® with four different Reynolds Averaged Navier 19 Stokes (RANS) turbulence model is constructed using a VOF model, to represent flows in this 20 manhole. Velocity profiles and pressure distributions from the models are compared with the 21 experimental data in view of finding the best modelling approach. It was found among four 22 different RANS models that the RNG k-ε and k-ω SST gave a better approximation for velocity and 23 pressure.

#### **Keywords** 24

25 Manhole, OpenFOAM<sup>®</sup>, RANS model, Stereoscopic PIV, VOF

#### Introduction 26

27 Manholes are one of the most common features in urban drainage networks. They are located at 28 changes in slope and orientation of the sewer pipes, as well as at regular intervals along the pipes 29 to enable maintenance. The flow pattern in a manhole is complex, especially during high flows, 30 and involves several hydraulic phenomena such as local flow contraction, expansion, rotation, 31 recirculation as well as possible air entrainment and sediment mixing. These flow phenomena can 32 control the overall energy loss, transport and dispersion of solute and particulate materials in the 33 manhole structure. PIV measurement can provide a good representation of the complex velocity 34 field of a manhole. Previously Lau (2008) studied two-dimensional PIV in a surcharged scaled 35 manhole. Attempts to measure stereo PIV data in a scaled manhole are however new and to the 36 authors' knowledge, has not been done before. Several researchers studied flow patterns in 37 surcharged manholes using CFD models. Use of different RANS modelling approach like RNG k-ε 38 model (Lau, Stovin, et al., 2007), Realizable k- $\epsilon$  (Stovin, Bennett, et al., 2013), k- $\omega$  model 39 (Djordjević, Saul, et al., 2013) have been reported. A little research study has been conducted on 40 how these flow patterns could affect flow movement and flow quality in the wider piped network. 41 In the current work, the flow phenomena of a scaled manhole are measured by stereo Particle 42 Image Velocimetry (PIV) and modelled numerically using OpenFOAM® CFD tools. Four different 43 RANS models, i.e. RNG k- $\epsilon$ , Realizable k- $\epsilon$ , k- $\omega$  SST and Launder-Reece-Rodi (LRR) were used, and 44 the differences in flow structures among them were compared and discussed.

### 45 Methods and Materials

### 46 **Experimental model**

47 The experimental facility was installed in the Hydraulic Laboratory of the University of Sheffield. It 48 consists of a transparent acrylic circular scaled manhole, linked to a model catchment surface 49 (Rubinato, Martins, et al., 2017). The manhole has an inner diameter ( $\Phi_m$ ) of 240 mm and 50 connected with a 75 mm diameter ( $\Phi_{\rm p}$ ) inlet-outlet pipes. Both pipes are co-axial, and the pipe 51 axis passes through the centre of the manhole vertical axis (Fig. 1). Two butterfly valves (one at 52 >48 $\Phi_p$  upstream the manhole, the other one at >87 $\Phi_p$  downstream the manhole) are used to 53 control the inflow and the water depth of the manhole respectively. The inflow was monitored 54 using an electromagnetic MAG flow meter fitted within the inlet sewer pipe ( $10\Phi_{p}$  from the 55 butterfly inlet valve). The ratio of the manhole diameters to inlet pipe diameters ( $\Phi_m/\Phi_p$ ) is 3.20. 56 Two Gems series pressure sensors (Product code 5000BGM7000G3000A, serial number 551362, 57 range 0-70 mb) were installed vertically at the inlet and outlet pipes, the first one at 350 mm 58 upstream from the centreline of the manhole, the second one, 520 mm downstream from the 59 centreline of the manhole by making a hole in the pipe of  $\Phi$  = 5 mm. They can measure 60 piezometric pressures for both free surface and pressure flow conditions within the inlet-outlet 61 pipes. These transducers were calibrated such that transducer output signal (4-20 mA) can be 62 directly related to gauge pressure. For the calibration of the pressure transmitter 10 different 63 water levels were measured (range 50 to 500 mm, no discharge) and checked with a point gauge. 64 Rubinato (2015) found an overall accuracy can be defined as ±0.72 mm. A transparent measuring 65 tape was attached vertically to the manhole side in order to check the manhole water levels 66 during the experiments. The tape position was at the other side of the camera, keeping an equal 67 distance from both inlet and outlet.



Fig. 1. Experimental setup for 2D3C stereo PIV measurement at the manhole

### 70 **PIV measurement**

68 69

71 A stereo PIV measurement setup was installed. Two Dantec FlowSense EO 2M cameras and an 72 Nd:YAG pulsed laser was placed at opposite sides of the scaled manhole. Each camera resolution 73 was 1600x1200 pixels (Fig. 1) and was set at the same distance from the manhole making more 74 than 45° angle at the vertical centre of the measuring plane. To reduce error due to refraction 75 through the curved manhole wall, a transparent acrylic tank was constructed around it and filled 76 with water, keeping flat surfaces parallel to both camera lens axes. The laser was directed from the bottom of the acrylic manhole as a laser sheet with the help of a flat mirror set at 45° to the 77 78 horizontal direction. The laser sheet thickness was around 4mm.

79 Conventional 2D PIV can give the velocity vectors perpendicular to the camera direction only, 80 which is typically parallel to the laser sheet (known as the in-plane velocity (Leandro, Bung, et al., 81 2014)). However, the use of two cameras can give the reading of the third component of the 82 velocity vector (referred as out of plane velocity) with proper regeneration of the 2D velocities 83 from each camera images; provided that appropriate calibration is done beforehand. As both of 84 the cameras see the same stationary image from two different angles, the calculated velocity 85 vectors from the two successive snapshots of a moving particle will also orient differently on both 86 cameras. This change in orientation is the result of the different camera position and out of the 87 plane velocity vectors. If the camera position is known and distortion of a still image at each 88 camera is known, the out of the plane velocity component can be calculated. Standard calibration 89 plates were used in this study to calibrate the cameras.

90 When the manhole surcharge level is below a certain limit (typically around  $0.2\Phi_m$ ), then the 91 manhole inlet flow reacts vigorously with the surface, creating irregular flow pattern and a very 92 high head loss. This surcharge limit is known as threshold surcharge (Stovin, Bennett, et al., 2013). 93 At higher surcharge, the flow pattern inside the manhole is regular. This work focuses to analyse 94 above threshold surcharge flow. The PIV measurements condition was chosen as 4 dm<sup>3</sup>/s of inflow 95 through the inlet pipe and a water level of 310 mm which resulted in a surcharge level of 235 mm 96 (s) at the manhole centre; making surcharge to manhole diameter ratio (s/ $\Phi_m$ ) 0.98. This was well 97 over the estimated threshold surcharge level. Initial inspection showed that the measurement 98 zone has two different distinctive velocity characteristics. One part of the measuring plane is 99 approximately in line with the inlet-outlet pipes and is characterised by a fast, slowly expanding jet 100 flow. The second part is outside the jet flow zone and is characterised by a recirculation in which 101 the velocity magnitude is around 10% of the jet flow. For these two distinctive velocity zones, data 102 were taken at different image time separation intervals ranging from 250 ms to 4000 ms; so that 103 velocities of both zones can be estimated accurately using the PIV cross-correlation algorithm.

For seeding, 100  $\mu$ m polyamide 12 particles were chosen (density = 1010 kg/m<sup>3</sup>). The particles were mixed with water and kept in a seeding tank with continuous circulation so that they remain in suspension. The particles were pumped from the seeding tank at a constant rate into the inlet pipe approximately 40 $\Phi_p$  upstream on the manhole so that they were well mixed before entering the manhole.

The seeding rate was adjusted by checking the PIV images to facilitate at least five particles in any selected interrogation area. Data were recorded at three vertical planes; one passing through the central axis of the inlet-outlet pipes and the other two at 50 mm offset from it (see Fig. 1). Each data set was measured for five minutes, at a rate of 8 image pairs per second, totalling 2400 pairs of images.

The data was analysed using Dantec Dynamics' DynamicStudio v3.31 software. The collected data was pre-processed after masking the area of interest. The fluid velocity was calculated using an adaptive cross-correlation technique keeping an interrogation area of 128x128 pixels with consideration of 50% overlap between two adjacent areas. Median correction post-processing was applied to remove erroneous vectors and which removed approximately 2 to 8% vectors from each measurement set.

Due to the resolution and positioning of the laboratory setup, neither of the cameras was able to cover the whole manhole height. Emphasis was given to the incoming jet to see how velocity is distributed over the length of the manhole. Hence the data was recorded covering the lower zone of the manhole; from the manhole bottom until the height of 150 mm of the manhole; which is two times of the inlet-outlet pipe diameter.

# 125 Numerical model

The open source CFD model tool OpenFOAM<sup>®</sup>v4.1 was used in this work. The solver *interFoam* is selected as it can predict the velocity patterns and the free-surface for sharp interfaces. This solver uses a single set of Navier-Stokes/Reynolds-Average equations where the velocity is shared by both phases and a Volume of Fluid (VOF) method (Hirt and Nichols, 1981) captures the free130 surface position. The length of the inlet pipe was chosen as 1000 mm (more than  $13\Phi_p$ ) based on 131 some other previous works (Lau, 2008; Stovin, Bennett, et al., 2013), and the outlet pipe was kept 132 as 400 mm, which is until the position of the pressure sensor at the downstream of the manhole 133 (Fig. 2a). The computational mesh for the simulation was prepared with hexahedral Cartesian 134 mesh using *cfMesh* (Juretić, 2015).

135 Some pre-analysis of CFD simulations were performed in order to test the mesh independence. 136 For this case, three computational meshes were constructed having dx = 2 mm (Mesh 1), 4 mm 137 (Mesh 2) and 6 mm (Mesh 3) respectively keeping the global refinement ratio as 3. Number of 138 cells at these meshes are: 2.4 million, 861,500 and 380,000 respectively. The inlet flow condition 139 was prescribed as constant discharge of  $Q = 4 \text{ dm}^3/\text{s}$ . The meshes were simulated using k- $\epsilon$ turbulence model and the velocity profiles at the manhole centre and at the outlet pipe were 140 141 extracted from the results (Fig. 2c). The mesh analysis was performed applying Richardson 142 extrapolation (Celik, Ghia, et al., 2008). The meshes gave similar results at the manhole jet zone 143 and at the pipe. However, the velocity profiles showed different results closed to the manhole 144 water surface (at around z = 0.29 m to 0.31 m) and close to the bottom (around z = 0 m). Mesh 3 145 predicted slightly slower velocity at the near surface zone. The apparent order (p) and the Grid 146 Convergence Index (GCI) was calculated at each grid point of the meshes. The average value of p 147 at the manhole centre and pipe were found 2.76 and 2.32 respectively. The GCI values were found 148 higher close to the surface and the walls. Analysis showed that 50% cells in the manhole has GCI 149 value below 10% when comparing Mesh 2 and Mesh 3; while 65% cell showed below 10% GCI in 150 case of comparing Mesh 1 and Mesh 2. However, in case of results at the pipe, 70% and 76% cells 151 showed GCI value below 10% in case of comparing Mesh 3-Mesh 2 and Mesh 2-Mesh 1 152 respectively. It was apparent that the results go towards mesh independence and Mesh 1 and 153 Mesh 2 show almost similar results. Average approximate relative error between Mesh 2 and 154 Mesh 1 was found to be 2.7%, whereas the simulation time requirement for Mesh 1 was more 155 than three times to that of Mesh 2. Considering the accuracy level and computational time 156 required, Mesh 2 with dx = 4 mm was found best suited for this work (y+ is around 5).

The model considers all the manhole borders as *noSlip* wall (i.e. zero velocity at the wall) and three open boundaries: *inlet, outlet* and *atmosphere* (Fig. 2b). Wall roughness was not considered as the aim was to characterize the flow velocity patterns in the manhole in which the wall energy losses were considered small in comparison with the entry, exit and mixing losses. The inlet boundary conditions were prescribed as fixed velocity approving for fully developed pipe flow profile using inverse power law of pipe flow (Çengel and Cimbala, 2006, chap. 8):

$$v_r = v_{max} \left( 1 - \frac{r}{R} \right)^{1/n} \tag{1}$$

163 where  $v_r$  is the longitudinal velocity at a radial distance of r from the pipe axis, R is the pipe 164 radius,  $v_{max}$  is the maximum longitudinal velocity at the developed profile section and n is a 165 constant which is dependent on the Reynold's number of the flow. To find the best combination of 166  $v_{max}$  and n, pre-analysis was done considering a pipe flow CFD model. The pipe diameter was 167 made the same as the inlet pipe ( $\Phi_p$ =0.075 m) and pipe length was kept 3 m (40  $\Phi_p$ ). The model 168 was simulated using k- $\epsilon$  model applying the same inflow (Q = 4 dm<sup>3</sup>/s). It was found that the pipe 169 becomes fully turbulent at flow reach of  $27\Phi_p$  (2 m length) and after 20 s of simulation time. It produces  $v_{max} = 1.128 \text{ms}^{-1}$  at fully turbulent condition and n = 6.5 gives the best fit curve of the development profile. These values were chosen to calculate inlet boundary condition of the manhole model using Eqn. 1. The outlet boundary condition was prescribed as fixed pressure boundary corresponding to average water column pressure head, measured with the outlet pipe pressure sensor (shown in Fig. 1). The pressure at atmosphere boundary condition (at the manhole top, shown at Fig. 2b) was prescribed as equal to *atmospheric pressure* and *zeroGradient* for velocity to have free air flow, if necessary.

177 The mentioned condition was simulated with four different Reynolds Average Navier-Stokes 178 (RANS) turbulence modelling approaches; namely: RNG k-ε model, Realizable k-ε model, k-ω SST 179 model and Launder-Reece-Rodi (LRR) model to evaluate if they are able to characterize the flow 180 properly. The first two models use a two-equation based approach calculating turbulent kinetic 181 energy (k) and turbulent energy dissipation ( $\varepsilon$ ). These two models are and formulated by Yakhot, Thangam, et al. (1992) and Shih, Liou, et al. (1995) and known to be better than the standard k- $\varepsilon$ 182 183 model to give better prediction at separating flow and spreading rate of round jets respectively. 184 The k- $\omega$  SST model used in OpenFOAM is based on Menter and Esch (2001) with updated 185 coefficients from Menter, Kuntz, et al. (2003) and addition of the optional F3 term for rough walls 186 (Hellsten, 1998). This model uses rate of dissipation ( $\omega$ ) instead of  $\varepsilon$  at the near wall zone and 187 standard k- $\varepsilon$  model at the zones far from the wall influence and is supposed to give better 188 prediction at the near wall and turbulence separating flow. Wall-functions are applied in this 189 implementation by using Kolmogorov-Prandtl expression for eddy viscosity (Hellsten, 1998) to 190 specify the near-wall omega as appropriate. The blending functions are not currently used in 191 OpenFOAM version because of the uncertainty in their origin. The effect is considered negligible in 192 case of small y+ cell at the wall (Greenshields, 2017) and hence can be applied to models with low 193 y+ cells. The fourth model uses the seven equation based Reynolds Stress Model (RSM); using 194 turbulent kinetic energy (k) and six component of stress tensor (R) directly and therefore may 195 predict complex interactions in turbulent flow fields in a better way. The initial condition was 196 prescribed as filling the manhole up to the expected level. Inlet pipe, outlet pipe and the manhole 197 zone in line with the inlet-outlet pipe was initialized with a fully developed velocity profile which 198 was same as the inlet boundary condition. The inlet turbulent boundary and initial conditions k,  $\varepsilon$ , 199 *R* and  $\nu_t$  were calculated using standard equations as follows:

$$k = \frac{3}{2} \left( I | \boldsymbol{u}_{ref} | \right)^2 \tag{2}$$

$$\epsilon = \frac{C_{\mu}^{0.75} k^{1.5}}{L_{1.5}} \tag{3}$$

$$v_t = C_\mu \frac{k^2}{\epsilon} \tag{4}$$

$$\omega = \frac{k^{0.3}}{C_{\mu}L} \tag{5}$$

where  $u_{ref}$  is the velocity to be considered, *I* is the turbulent intensity (chosen as 0.05 at this case),  $C_{\mu}$  is a constant (=0.09) and *L* is the characteristic length of the inlet pipe.

202 Standard wall function was considered in k- $\omega$  SST model only as according to the model 203 description, it requires wall function for y+>1. The rest of the models require wall function when 204 30 < y < 300.



205

206Fig. 2. a) Numerical model mesh (top left panel), b) Boundary locations (top right panel) and207c) Mesh convergence test (bottom panel)

During the simulations, the *adjustableRunTime* code was used keeping maximum Courant– Friedrichs–Lewy (CFL) number to 0.95. Cluster computing system at the University of Coimbra was used to run the simulations using MPI mode. Each simulation was run for 65 seconds. The steady condition was ensured by checking the residuals of p, alpha.water, k, epsilon and Rxx (where applicable). The first 60 seconds were required to reach steady state condition and the results of 101 velocity profiles in the last 5 seconds were averaged and all the numerical analysis were made using averaged data of these mentioned time step results.

# 215 **Results and Discussion**

216 The processed PIV velocity data were compared with velocity data of the manhole model. The PIV 217 measurement was taken at the central vertical plane (CVP) along with the left vertical plane (LVP) 218 and right vertical plane (RVP) (Fig. 1). However, the position of LVP was much further from the 219 camera, and light ray coming from the two edges of the LVP have to travel through the manhole 220 inlet-outlet pipe. Although the scaled manhole model is made of transparent acrylic, the joints 221 between the model components were semi-transparent to non-transparent. Due to this limitation, 222 the LVP image edges had bad data in some cases, as the camera could not see the seeding 223 particles due to the obstruction made by the model joints. So in this work, PIV data comparison is only done to the CVP and RVP of the manhole. Fig. 3 shows the axial  $(V_x)$  and vertical  $(V_z)$  velocity components comparison between the PIV data and numerical model results. All the comparison is shown as dimensionless velocities as a ratio of the average inlet velocity  $(V_{avg})$ , where  $V_{avg}$  is the ratio between inlet discharge (Q) and pipe cross section  $(A_p)$ . Both PIV data and CFD data are showing temporal mean velocities from the measurements and simulation results. PIV data at the CVP near the manhole inlet and outlet pipe were not collected as it was not visible clearly by the cameras due to the joints between pipe and manhole.

It can be seen from the velocity comparison at the CVP of PIV (1<sup>st</sup> and 3<sup>rd</sup> column of row 1 at Fig. 231 3), that the jet flow in the experimental results starts expanding slowly as it proceeds from the 232 inlet towards the outlet. The CVP axial velocity  $(V_x)$  (1<sup>st</sup> column of Fig. 3) reaches up to 110% of  $V_{avg}$ 233 at the experimental results. It dampens down at the manhole centre. The maximum velocity near 234 235 the outlet pipe is the same as V<sub>avg</sub>. However, in the CFD data, this damping effect is not seen. The 236 high-velocity core stays as 110% of V<sub>avg</sub> until it reaches the outlet. While the jet flow trying to 237 escape through the outlet pipe, it hits the manhole wall at the top of the outlet (near x = -100 mm 238 and y = 75 mm) and creates a vertically upward velocity component which is up to 20% of  $V_{avg}$  (3<sup>rd</sup> column of Fig. 3). However, in all the CFD, the vertical velocity component at this zone reaches up 239 240 to 30% of V<sub>avg</sub>. Comparing the velocity contours of V<sub>x</sub> and V<sub>z</sub> at the CVP, it can be seen that LRR 241 model could not predict the velocity profile properly. This model shows different axial velocity

242 contour at CVP which is dissimilar to PIV measurements.



243

Fig. 3. Comparison of non-dimensional velocity components from the numerical models and PIV
 measurement at both CVP and RVP of the manhole. The flow direction is from right to the left

Comparing the velocity towards the pipe axis at the RVP ( $2^{nd}$  column of Fig. 3), RNG k- $\varepsilon$  and 246 Realizable k-ɛ models slightly underestimate the axial velocity component in compared to the PIV 247 experiment. The PIV data showed axial velocity up to 0.3Vavg whereas, in these two numerical 248 models, the highest velocity is found 0.25V<sub>avg</sub>. The axial velocity at this plane is properly estimated 249 by k-ω SST model. The shape of the velocity contour is almost similar to that of the PIV data. The 250 vertical velocity component ( $V_z$ ) (4<sup>th</sup> column of Fig. 3), at PIV is observed between -0.2 $V_{avg}$  and 251 0.2V<sub>avg</sub>. The numerical models show similar results of V<sub>z</sub> near the outlet of the manhole. However, 252 at the upper part of the measuring plane near the inlet, the Vz was measured in PIV as around -253  $0.2V_{avg}$ , which was not predicted by the numerical models. Only k- $\omega$  SST model in this plane shows 254 255 a negative (downward) velocity near the inlet, however still underestimated than the PIV.

The velocity comparison was not done for out of the plane component  $(V_y)$  as at CVP, this component is very close to zero and considered not significant (Fig. 4).



# 258 259

Fig. 4. Out of the plane velocity component at the CVP from PIV and four RANS models

The standard deviation of the data was also compared at both CVP and RVP for axial and vertical
 velocity components (Fig. 5). The comparison was also made on dimensionless velocities as a ratio
 of V<sub>avg</sub>.

It can be seen that for axial velocity component ( $V_x$ ) at CVP and RVP (1<sup>st</sup> and 2<sup>nd</sup> column at Fig. 5), 263 the PIV data shows high standard deviation near inlet, outlet and jet expansion zones. High 264 265 standard deviation can also be seen near the manhole floor of CVP. Vertical velocity component  $(V_z)$  showed lower deviation compared to  $V_x$ . The standard deviation was found to be significantly 266 low at both planes in CFD results in compared to PIV data. All CFD results show marginally higher 267 268 standard deviation values close to the jet expansion zone. The Realizable k-E model shows the 269 minimum fluctuation in the velocity fields, resulting almost zero standard deviation. The k- $\omega$  SST 270 model shows the highest standard deviation among all the four RANS models, still, the value is 271 lower than that of the PIV data. As these numerical models are formulated from Reynolds 272 Averaged Navier Stokes (RANS) equations, a big part of the turbulence variabilities is averaged out 273 from the results already. Perhaps Large Eddy Simulation (LES) model could show better standard 274 deviation match with the PIV data. Moreover, for this work, only 5 seconds of numerical 275 simulation data was considered. In case the longer period of data was taken into consideration, 276 probably the standard deviation in the CFD results would come higher. Turbulence coming from 277 pipe rather than the standard turbulence inlet conditions could also be a reason for the 278 differences in numeric.

Examining the velocity contours at Fig. 3 and Fig. 5, it can be stated that  $k-\omega$  SST model creates the closest approximation of the manhole velocity field followed by RNG k- $\epsilon$  model. The LRR model overestimates the axial velocity component in this plane and the velocity contour is significantly different than the PIV measurement.



283

284 285

Fig. 5. Temporal standard deviation of different velocity component at CVP and RVP, measured from PIV data and four RANS models

To understand the manhole flow more clearly, streamline and water level fluctuation was 286 287 analysed from the numerical model results (Fig. 6). It can be seen that the inflow jet originating 288 from the inlet, marginally spread in the manhole (termed as the diffusive region). When exiting 289 the manhole through the outlet pipe, the diffusive region impinges the manhole wall and after 290 that part of the jet flow moves vertically upward. This flow region reaches the manhole water 291 surface, starts moving opposite to the core jet direction and makes a clockwise circulation. Due to 292 this constant circulation, the free surface of the manhole fluctuates slightly. The fluctuation was 293 found more towards the inlet of the manhole. In this case, the water level was observed varying 294 between 0.295 m to 0.320 m. The fluctuation level was found different in the four numerical 295 models. In RNG k-ε model, the water level varies from 0.300 m to 0.310 m. While the water level 296 fluctuations in the other three manhole models were found as Realizable k-E: varying between 297 0.300 m and 0.305 m; k-ω SST: varying between 0.300 m and 0.310 m and LRR: varying between 298 0.310 m and 0.320 m. In the experimental work the average water level at the manhole was 299 observed as 0.310 m. The fluctuation of the manhole level was not possible to measure in the

- 300 experimental work as it would require installing a pressure sensor at the manhole bottom, which
- 301 would create an obstacle in the laser ray path line and hence was not used.



302303

Fig. 6. Flow streamline through the manhole and water level range from different models

304 The pressure distributions in different CFD models were also compared with experimental data 305 and can be seen in Fig. 7. All distances showed at the horizontal axis are measured from the 306 manhole centre. The inlet and outlet pipes are connected at distance of 0.12 m and -0.12 m 307 respectively. The bottom pressure in between these two distances also represents free surface 308 water level inside the manhole. Two box plots represent pressure data recorded during the 309 experimental measurement using pressure sensors. The left box plot shows the pressure at the 310 outlet pipe, whose average value was used to generate boundary condition of the numerical 311 models. As the downstream pressure is same for all the four models, all the model results pass 312 through the average value of this box. Pressures at the upstream side were calculated from the 313 models. All the line plots are showing maximum, minimum and average bottom pressure from 314 each of the four CFD models. The circular marker at the manhole centre (at x=0 m) is showing the 315 product of recorded average water height at the manhole during the experimental works (h = 316 0.310 m), water density ( $\rho$ ) and gravitational acceleration (g).



#### 317

Fig. 7. Bottom pressure comparison from different RANS models and the experimental pressure
 sensor data

320 From Fig. 7, it can be seen that when the upstream flow through the inlet pipe enters the 321 manhole, the flow experiences a pressure drop, which is due to the expansion of the flow. The 322 second pressure drop can be observed when the flow exits the manhole and enters the outlet 323 pipe. This drop is due to the flow contraction and much bigger in magnitude. The average pressure 324 line of RNG k- $\varepsilon$ , Realizable k- $\varepsilon$  and k- $\omega$  SST models produce a similar pressure pattern throughout 325 the computational domain. The LRR model overestimates the bottom pressure of the manhole, 326 although the difference found is in the range of few millimetres of water column head. The 327 maximum, minimum and the average bottom pressure from RNG k-E model shows almost the 328 same line which presents that the RNG k- $\varepsilon$  model shows almost zero pressure fluctuation. 329 Apparently, the LRR model overestimated the pipe loss in compared to the other three models. 330 This results in an overestimation of upstream pressure at the inlet direction.

The coefficient of head loss (K) in the manhole is known as the ratio between head loss and the velocity head and is calculated using equation (6).

$$K = \Delta H / \left(\frac{\nu^2}{2g}\right) \tag{6}$$

where  $\Delta H$  is the head loss, v is the average longitudinal velocity at the outlet pipe (= 0.89 m/s) and g is the acceleration due to gravity.

As the numerical model reached steady state before extracting any results, and as both inlet and outlet pipes were full, the temporal averaged velocity at each pipe can be considered equal. In this case, a difference in bottom pressure would give the same value as head loss. To compute the pressure drop at the manhole centre for a certain CFD result, each line showing the average bottom pressure (Fig. 7) was projected to the manhole centre from both inlet and outlet pipe. The 340 vertical difference of pressure value between these two lines at the manhole centre gives the

341 value of pressure drop for the manhole, which is later divided by  $\rho g$  and considered as  $\Delta H$ .

342 The value of manhole head loss coefficient has been reported in different literature. This is directly 343 related to the structural mould types of the manhole, manhole to pipe diameter ratios, as well as 344 manhole surcharge ratio. A comparable analysis of the coefficients of head loss (K) with those of 345 the values reported in different literature are shown in Tab. 1. It should be noted that only the 346 research works reporting the same manhole mould type are considered here. In the work of both 347 Arao and Kusuda (1999) and Lau, Stovin, et al. (2008), authors reported high head loss coefficient 348 at below threshold surcharge conditions and comparably lower coefficient at above threshold 349 surcharge condition. As in this research, the manhole surcharge condition is comparable to above 350 threshold surcharge, only the coefficient range covering this condition are shown.

Works done	Head loss coefficient	Experimental condition $(\Phi_m/\Phi_p)$	Surcharge ratio range	
			s/Φ <sub>m</sub>	s/Φ <sub>p</sub>
Marsalek (1981)	0.210	1.923		
Arao and Kusuda (1999)	0.18-0.58	3.60	0.55-1.67	2-6
Lau, Stovin, et al. (2008)	0.28-0.69	9.08	0.65-0.82	5.90-7.45
This work				
RNG k-ε	0.193	3.20	0.98	3.13
Realizable k-E	0.156			
k-ω SST	0.284			
LRR	0.265			

351 Tab. 1. Different values of head loss coefficient (K) at different models

### 352

353 Tab. 1 shows that the four models calculate the manhole flow differently and hence give different 354 values of manhole head loss coefficient. It is reported by different authors the head loss 355 coefficient becomes higher when manhole to pipe diameter ratio  $(\Phi_m/\Phi_p)$  is high and vice versa 356 (Stovin, Bennett, et al., 2013; Bo Pedersen and Mark, 1990). The manhole reported at Arao and 357 Kusuda (1999) has a similar  $\Phi_m/\Phi_p$ . Comparing the findings from the literature, it is apparent the 358 Realizable k- $\varepsilon$  model gives rather low head loss coefficient for this case. From the rest three 359 models, the coefficient given by k- $\omega$  SST model is almost 50% more than that of value given by 360 RNG k- $\varepsilon$ , however, both values lie within the range specified by other researchers.

# 361 Conclusions

362 In this work, two-dimensional three component (2D3C) stereo PIV measurement was done on a 363 scaled inline manhole with manhole to pipe diameter ratio of 3.20, in order to evaluate CFD model 364 constructed in OpenFOAM<sup>®</sup> and four different RANS models with VOF method. From the analysis, 365 it can be apparent that each model calculates the velocity inside manhole differently. Comparison 366 with PIV measurement at the central vertical plane showed similar velocity as compared to all the 367 numerical models. However, comparison of the velocity at another vertical plane 50 mm offset to 368 the centre, showed that all the CFD models slightly underpredicts the axial velocity. The velocity 369 and locations of vortex structures centres were found marginally different among the models. The

370 k-w SST model showed the closest approximation of velocity contour followed by the RNG k-e 371 model. All the models showed very good approximations of the average water surface level at the 372 manhole. However, the LRR model could not quite capture the velocity profile in compared to PIV 373 data. This model predicts considerably higher axial velocities for central vertical plane as 374 compared to PIV. Nonetheless, the temporal standard deviations of the axial and vertical velocity 375 components were found significantly low when compared to those of experimental measurement 376 through PIV. As a RANS model is formulated based on time-averaged turbulence data, which could 377 be the reason of having a lower standard deviation in the CFD model.

Bottom pressure analysis through the computational domain shows that the average pressure line is almost similar at RNG k- $\varepsilon$ , Realizable k- $\varepsilon$  and k- $\omega$  SST models. However, the comparison could be made with data from only two pressure sensors installed at the inlet and outlet pipe respectively. RNG k- $\varepsilon$  model showed almost no pressure fluctuation while maximum pressure line predicted by the Realizable k- $\varepsilon$  model was found much higher than the measurement. The calculated head loss coefficients were compared with the values reported in the literature. It was seen that the Realizable k- $\varepsilon$  model shows much lower value compared to the values reported.

385 Considering all the aspects of the four models analysed here, it can be said that both RNG k- $\epsilon$ 386 model and k- $\omega$  SST models give a very good approximation of manhole hydraulics. However, it 387 should be noted while using k- $\omega$  SST model, the wall boundary cell size must be made considerably 388 small for the proper formulation of the model.

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